Solid State Communications, Vol. 9, pp. 1057-1060, 1971. Pergamon Press.

Printed in Great Britain

## LOW-SPIN → HIGH-SPIN STATE TRANSITION IN HIGH PRESSURE COBALT SESQUIOXIDE

J. Chenavas and J.C. Joubert

Laboratoire d'Electrostatique et de Physique du Metal C.N.R.S., Cedex No. 166, 38-Grenoble Gare, France

and

## M. Marezio

Bell Telephone Laboratoires, Incorporated, Murray Hill, New Jersey 07974

(Received 28 April 1971 by E.F. Bertaut)

 $\text{Co}_2\text{O}_3$  with the corundum structure was synthesized under high pressure. A refinement of the structure indicated that  $\text{Co}^{3+}$  is in the low-spin state. By annealing, this oxide transformed to a corundum structure with a larger volume indicating that the  $\text{Co}^{3+}$  is in the high spin state. The increase in volume was 6.7 per cent.

ALL TRIVALENT 3d transition metals form a sesquioxide with the corundum structure except cobalt, nickel and manganese. The former two elements are more stable as divalent cations. CoO and NiO with the NaCl structure are the stable oxides. Another cobalt oxide is known, namely Co3O4, which crystallizes with the normal spinel structure. In this oxide 2/3 of the cobalt ions are trivalent while the other third is divalent. It has been shown that the Co3+ ions in this compound are in the low spin state. Only a few cobalt compounds are known in which the cobalt ions are all trivalent, LiCoO2 and CoF3 are two examples. From the ionic radii of Co3+ and Co2+ with coordination number 6, one could speculate that the transformation Co2+ → Co3+ should be favored by high pressures, especially if the Co3+ ions are in the low spin state (see Table 1). This note reports the high pressure synthesis of Co<sub>2</sub>O<sub>3</sub> with the corundum structure.

Two different syntheses were performed. In the first the starting material was  $Co^{2+}Cr^{6+}O_4$ . This double oxide has the orthorhombic

Table 1

Valence	Coord.	Spin	Radius (Å)
3	6	Low	0.525
2	4	High	0.57
3	6	High	0.61
2	6	Low	0.65
2	6	High	0.74

Cr³+V⁵+O₄ structure in which the Co²+ ions have octahedral coordination, while the smaller Cr⁵+ ions are tetrahedrally coordinated. The high pressure runs were done in a modified 'belt' aparatus. A platinum cylindrical cell was used as sample container. A standard heater assembly was used. The pressure was calibrated by measuring the resistance of bismuth and thallium wires at room temperature. The temperature was calibrated by a Pt-Pt/10% Rh thermocouple. By applying a pressure of ≥ 90 kbar and 1000°C in the presence of a small amount of water, CoCrO₄ decomposes into Co₂O₃ with the corundum structure and a solid solution of CoOOH and CrO₂ with the

T

aı

a

S

T

Iı

orthorhombic distortion of the rutile structure, as revealed by X-ray photographs of the quenched materials. The decomposition occurs according to the formula:

$$3 \text{ CoCrO}_4 + \frac{1}{2} \text{ H}_2 \text{O} \rightarrow 4 \text{ (CrO}_2)_{0.75} \text{(CoOOH)}_{0.25} + \text{Co}_2 \text{O}_3 + \frac{3}{4} \text{O}_2$$
 (1)

Under pressure  $\text{Co}_2\text{O}_3$  transforms in the presence of water to a high pressure form of CoOOH, therefore different amounts of water affect the composition of the solid solution  $(\text{CrO}_2)_x(\text{CoOOH})_{1-x}^3$ 

The second synthesis was based on the following reaction:

$$2 \text{ CoF}_3 + 3 \text{ Na}_2 \text{O}_2 \rightarrow 6 \text{ NaF} + \text{Co}_2 \text{O}_3 + \frac{3}{2} \text{ O}_2$$
 (2)

The X-ray pattern of the quenched material of a run done at 80 kb and 850°C showed the same hexagonal pattern obtained in syntheses.(1). Unfortunately some reduction took place and the quenched material contained some Co<sub>3</sub>O<sub>4</sub>.

Table 2. Powder pattern of Co2O3(low spin)

hkl	$d_{0}$	$d_c$	$I_{o}$
012	3.476	3.489	25
104	2.546	2.551	100
110	2.389	2.391	66
113	2.088	2.092	16
024	1.743	1.745	33
116	1.601	1.602	90
214	1.408	1.409	43
027	1.379	1.380	50
1010	1.236	1.236	48
119	1.233	1.233	15
220			
217	1.195	1.195	46
306	1.163	1.163	48

Table 2 gives the powder pattern of  $\text{Co}_2\text{O}_3$  which was indexed on a hexagonal cell,  $a = 4.782\,\text{Å}$ ,  $c = 12.96\,\text{Å}$ , c/a = 2.710,  $V = 258.5\,\text{Å}$ . These values correspond to the material of synthesis (2). The X-ray powder photographs were taken using a Norelco camera of

114.6 mm dia. and filtered CrKa radiation. In general the synthesis (1) yielded lattice parameters which were slightly larger. A typical run gave  $a = 4.790 \,\text{Å}$  and  $c = 13.00 \,\text{Å}$ . We believe that these larger values are due to a small percentage of substituted Cr. The pattern reported in Table 2 resembles closely that of  $\alpha - Al_2O_3$  as far as the sequence of the d-spacings is concerned, but the intensities of the reflections are similar to those of V2O3. This indicates that the radius of the Co3+ ions is very close to that of the Al3+ ions. Therefore the Co3+ ions must be in the low spin state. The reported ionic radius for the six coordinated Al3+ is 0.53 Å, while that of the corresponding Co<sup>3+</sup> in the low spin state is 0.52 Å. On the other hand the arrangement of the structure of Co<sub>2</sub>O<sub>3</sub> should be different from that of Al<sub>2</sub>O<sub>3</sub> as the relative intensities indicate.

By annealing Co<sub>2</sub>O<sub>3</sub> at 400°C for a half hour in air, this oxide transforms to a corundum structure with a different volume and c/a ratio. The new values are  $a = 4.882 \,\text{Å}$ ,  $c = 13.38 \,\text{Å}$ , c/a = 2.742,  $V = 276.2 \text{ Å.}^3$  The X-ray powder pattern is reported in Table 3. Since the unit cell volume and the d-spacings of this Co2O3 are nearly the same as those of a-Fe<sub>2</sub>O<sub>3</sub>, the Co3+ ions must be in the high spin state. The oxide Co2O3 with Co3+ ions in the low spin state will be referred herein as Co2O3 (low spin) whereas that with Co3+ ions in the high spin state will be referred as Co2O3 (high spin). The transformation Co<sub>2</sub>O<sub>3</sub> (low spin) → Co<sub>2</sub>O<sub>3</sub> (high spin) is accompanied by an increase in volume  $\Delta V/V(\text{low spin}) = 6.7 \text{ per cent. These}$ results seem to indicate that Co2O3 (high spin) is the cobalt sesquioxide stable at atmospheric pressure whereas Co2O3 (low spin) is the high pressure phase. This pressure-induced transformation involves only a change in the spin state of the cations. The coordination of the ions and their packing remain the same. The coordination-number of the cations remains six and for the anions it remains four, while the oxygen packing is nearly h.c.p. in both oxides.

Preliminary refinements of the two structures based on powder intensities gave the positional parameters and the interatomic distances reported in Table 4. The average Co-O distances are 1.90 n. In paracal run

to a pattern at of

ies of O<sub>3</sub>.

+ ions
nerefore
ate.
rdinated
onding
the
re of

 $1_2O_3$  as

half undum ratio. 3.38 Å, owder unit o<sub>2</sub>O<sub>3</sub>, the The spin ow high h spin). Co<sub>2</sub>O<sub>3</sub> in

high
h spin).
Co<sub>2</sub>O<sub>3</sub>
in
These
n spin)
pheric
high
anspin
the
Che
s six
he

ructures itional reported are 1.90

oxides.

Table 3. Powder pattern of Co2O3 (high spin)

hkl	$d_{0}$	do	Io
012	3.571	3.574	24
104	2.620	2.624	61
110	2.440	2.441	50
113	2.139	2.141	21
024	1.788	1.787	24
116	1.647	1.647	100

Due to the short annealing time, the back-reflection lines were diffuse. Extended annealing times produced the reduction of  $Co_2O_3$  to  $Co_3O_4$ .

and 1.99Å for Co2O3 (low spin) and (high spin) respectively. These values are in fairly good agreement with those calculated from Shannon and Prewitt ionic radii, 1.92 Å and 2.01 Å respectively. The oxygen octahedra in the two compounds seem to have different distortions. That of the Co2O3 (high spin) is similar to the one found in the other corundum structures. Instead, the distortion of the Co<sub>2</sub>O<sub>3</sub>(low spin) seems quite unique. The Co-O distance toward the shared face is shorter than that toward the unshared face. The contrary is true for all the other corundum structures. The Co-Co distances across the shared face between two octahedra are quite short in both compounds. Relative to the ionic radii it is smaller in Co2O3(high spin). Also the Co-Co distance across the shared edge is anomalously short in Co2O3 (high spin).

Due to the contamination of the samples from the by-products of both reactions we have been unable to measure physical properties such as resistivity and magnetic susceptibility.

In the last decade the transition metal sesquioxides with the corundum structure have been thoroughly studied because of their quite unique electrical and magnetic properties. As the number of d-electrons of the ions,  $n_d$ , increases the transition metal oxides go from a band metal behavior to a localized insulator behavior.  $^{4}$  Also the oxides with  $n_{d} \geqslant 2$  are magnetically ordered at low temperatures whereas those with  $n_d < 2$ , such as  $Ti_2O_3$ , have not been found to order at any temperature. The oxides with a large number of d-electrons, such as Cr2O3 and Fe2O3 are magnetic insulators. According to these general trends, Co2O3 (high spin) should be a magnetic insulator, whereas Co2O3 (low spin) a nonmagnetic insulator. In octahedral coordination the Co3+ ions in the low spin state have the configuration  $t_{20}^6 e_a^0$ which corresponds to a zero magnetic moment. The filled  $t_{2g}$  band would be compatible with the predicted insulator character of Co2O3(low spin).5

From shock-wave experiments it has been shown that Fe<sub>2</sub>O<sub>3</sub> (high spin) undergoes a phase transition from corundum to a denser phase.<sup>6</sup>
The data extrapolate to a zero-pressure density of 5.96 g/cm<sup>3</sup> for the new phase, as compared

Table 4. Positional parameters

		Low Spin	High Spin
Zco		0.346 ± 0.002	$0.346 \pm 0.002$
x <sub>o</sub>		$0.295 \pm 0.013$	$0.344 \pm 0.020$
$R = \Sigma \Delta F / \Sigma F$		0.03	0.04
,	Interatomic distanc	es in Co <sub>2</sub> O <sub>3</sub>	
Co-O	toward shared face	$1.88 \pm 0.06$	$2.12 \pm 0.08$
Co-O	toward unshared face	$1.29 \pm 0.04$	$1.86 \pm 0.04$
0-0	shared edge	$2.60 \pm 0.02$	$2.79 \pm 0.07$
0-0	shared face	$2.44 \pm 0.11$	$2.91 \pm 0.17$
0-0	unshared edge	$2.74 \pm 0.02$	$2.75 \pm 0.03$
0-0	unshared face	$2.93 \pm 0.06$	$2.77 \pm 0.08$
Co-Co	across face	$2.49 \pm 0.05$	$2.58 \pm 0.06$
Co-Co	across edge	$2.78 \pm 0.06$	$2.84 \pm 0.08$

with the value of  $5.24\,\mathrm{g/cm^3}$  observed for  $\mathrm{Fe_2O_3}$  with the corundum structure at zero pressure. Several explanations based on crystal chemical relationships have been offered for this transformation. The high spin-low spin transition found in  $\mathrm{Co_2O_3}$  offers another plausible one. The  $\mathrm{Fe^{3+}}$  ion in the low spin state has an ionic radius of 0.55 Å. From this value one can interpolate the molecular volume and the density

at zero pressure of  ${\rm Fe}_2{\rm O}_3({\rm low\ spin})$ . The calculated density of this phase, 5.98  $\pm$  0.06 g/cm³, seems to indicate that the transition observed by shock-wave in  ${\rm Fe}_2{\rm O}_3$  is  ${\rm Fe}_2{\rm O}_3$  (high spin)  $\rightarrow$   ${\rm Fe}_2{\rm O}_3$  (low spin).

Acknowledgements — The authors would like to thank P.D. Dernier and J.J. Capponi for valuable discussions.

## REFERENCES

- All the ionic radii quoted in this paper are from: SHANNON R.D. and PREWITT C.T., Acta Crystallogr. B25, 946 (1969), B26, 1046 (1970).
- 2. JEFFERY R.N., BARNETT J.D., VANFLEET H.B. and HALL H.T., J. appl. Phys. 37, 3172 (1966).
- 3. CHENAVAS J., JOUBERT J.C. and MAREZIO M., to be published.
- 4. McWHAN D.B., MENTH A. and REMEIKA J.P., J. Phys. (in press).
- 5. RICE T.M., (private communication).
- 6. CLARK S.P., Ed. Handbook of Physical Constants, Rev. Ed., Geol. Soc. Am. Mem. 97, 153 (1966).
- 7. REID A.F., and RINGWOOD A.E., J. Geophys. Res. 74, 3238 (1969).
- 8. AHRENS T.J., ANDERSON D.L. and RINGWOOD A.E., Rev. Geophys. 7, 667 (1969).

L'oxyde de cobalt  $\mathrm{Co_2O_3}$  a été synthétisé sous haute pression. Un affinement de la structure montre que l'ion  $\mathrm{Co^{3^+}}$  se trouve dans l'état 'low-spin'. Après recuit, cet oxyde se transforme en une nouvelle phase de structure corindon et de densité plus faible, contenant l'ion  $\mathrm{Co^{3^+}}$  dans l'état 'high spin'. L'augmentation de volume correspondant est de 6.7 pour cent.